



k-dependence of the mass renormalisation in $\text{Bi}_2\text{Sr}_2\text{CaCu}_2\text{O}_{8+\delta}$: a polarisation-dependent ARPES study



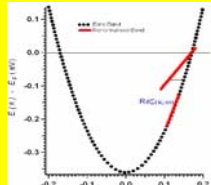
I.Santoso¹, S. de Jong¹, F.Massee¹, Y. Huang¹, A. Mans¹, W. K. Siu¹, R. Follath², M. Shi³, L.Patthey³, M.S. Golden¹

1. Condensed Matter Physics Group, Van der Waals-Zeeman Institute, University of Amsterdam, Netherlands,
2. BESSY GmbH, Einsteinstrasse 15, D-12489 Berlin, Germany. 3. Swiss Light Source, CH-5232, Villigen PSI, Switzerland

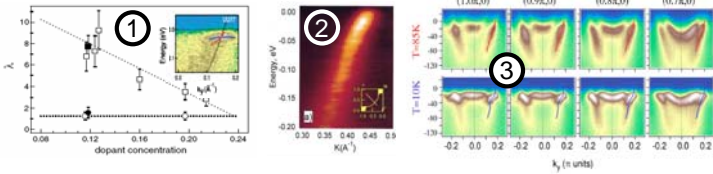


1. Introduction

We have used polarised Angle Resolved Photoemission (ARPES) to study the mass renormalisation around the Fermi surface in Pb-doped $\text{Bi}_2\text{Sr}_2\text{CaCu}_2\text{O}_8$. This renormalisation is seen as a deviation of the $E(k)$ dispersion of the independent particle band at low energy, close to the Fermi level.



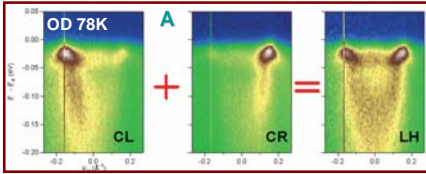
Mass renormalisation: due to “dressing” of electron with interactions i.e. a coupling to collective, bosonic modes



- In Bi2212 , the $(\pi,0)$ (=antinode point) renormalisation has a strong doping dependence Kim *et al.*, PRL91,167002 (2003).
- It shows doping dependences in the nodal direction (Bogdanov *et al.*, Valla *et al.*)
- In OD 71, a momentum and temperature dependence was found near $(\pi,0)$. These measurements were from a sample with superstructure modulations Gromko *et al.*, PRB68, 174520 (2003).

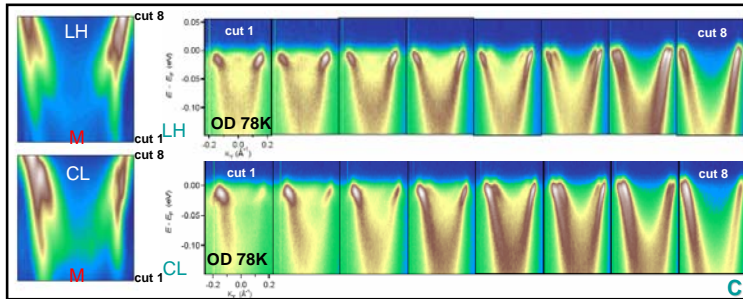
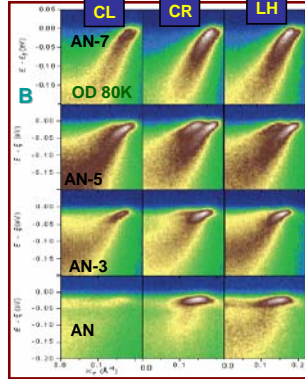
- What is the k -dependence of this renormalisation ?
- Exploit variable polarisation to help distinguish the BB and AB bands

4. Results and data analysis OD 78K & 80K: polarisation-dependence



Panel (A): CL and CR each pick out one side of the $(0,0)$ - $(\pi,0)$ line. The LH Spectra is very close to the sum of CL + CR.

Panel (B): zoom of positive k_y side of the spectra for LH, CR, CL polarisations from AN to 7° towards the node. Trend : CR favours AB band, CL the BB. LH is in between.

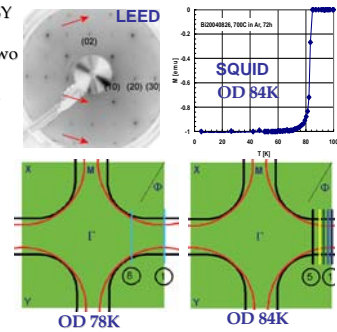


Panel (C): Momentum dependence of renormalisation for LH and CL polarised radiation. The degree of renormalisation is maximal at the AN and decreases as one moves away from the AN.

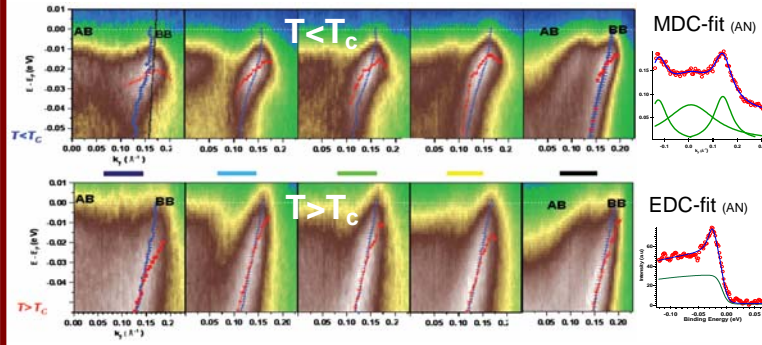
Conclusions from panels B and C: exploiting circular light enables better resolution of the c-axis bilayer splitting. Multiple EDM's per cut make MDC and EDC fitting more robust. This polarisation dependent data is still under analysis.

2. Experimental

- SIS beamline X09LA at SLS and U125-1/PGM at BESSY
- SFS2002 (SLS) and SFS-100 (IFW Dresden@BESSY)
- hv chosen as 39 eV: maximises contrast between the two c-axis split bilayer bands
- (Pb,Bi)-2212 single crystals (FZ) grown in Amsterdam Untwinned, modulation-free.
- Tc of these crystals: OD 78K, OD 80K, and OD 84K
- For OD 78 and 80K, data taken from 8 different cuts and using different polarisations:
 - linear horizontal (LH)
 - circular left (CL)
 - circular right (CR)
- Cut 1 represents antinode (AN), cut 8 is 7° away
- For OD 84K, 5 cuts are taken using LH radiation. Cut 1 = AN and cut 5 is 6° away from AN.



3. Results and data analysis: OD 84K



- Peak position in EDM shows band dispersion. EDC (red) dispersion more sensitive to the presence of the gap compared to the MDC (blue) dispersion.
- MDC dispersion is fitted using simple Lorentzian while EDC is fitted using Lorentzian multiplied by the Fermi-Dirac distribution, including a scaled background from $k > k_F$.
- Data clearly resolve the bilayer splitting

EDC (red) dispersion can be approximated by a Bogolubov quasiparticle dispersion in the presence of the gap:

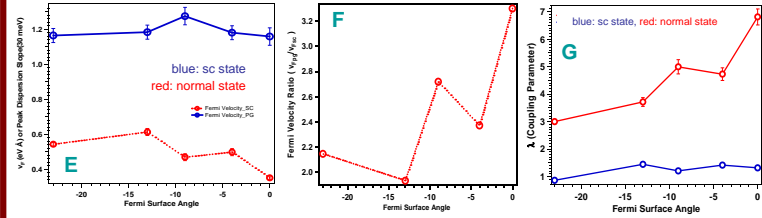
$$E(k) = \sqrt{m(k - k_F)^2 + \Delta^2} \quad \text{where} \quad m = \frac{v_F^2}{(1 + \lambda)} \quad (1)$$

Eq.1 works well at weak coupling regime. At strong coupling regime, the energy gap is similar to mode energy. It needs more complex expression (Fink *et al* CM/0512307)

The bare Fermi velocity can be obtained from the LDA data for the separation between bottom of the bonding (BB) and antibonding (AB) bands, given the experimental k_F for the BB:

$$v_F^0(k) = \frac{E_{\text{bond}}(k^2 - k_F^2)}{k_F^2} \quad (2)$$

- The slope of (2) at low energy is bare velocity
- BCS-like back dispersion in the superconducting state (particle-hole mixing), see Fig. H



- Trend: for T above Tc, the slope is constant (E).
- Trend: degree of renormalisation (seen as the slope change) in the superconducting state decreases as we go from the antinode half-way towards the node (E and F)
- The coupling constant change on leaving the sc state decreases as we move from antinode towards the nodal line (G), while it is more like constant at T above Tc (G).

5. Conclusions

- The mass renormalisation for the bonding band for $T < T_c$ is significantly momentum dependent:
 - It is maximal at the antinodal point, decreasing steadily as one goes towards the node. This k -dependence practically vanishes for T above Tc.
- Use of circular polarised radiation offers bright prospects for improved BLS contrast in the tricky zone between the antinode and node.

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Contact: Iman Santoso isantoso@science.uva.nl